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Magnetic Field Effects on Endcap EM Calorimetry in SDC

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1. Introduction

The SDC calorimeter will be immersed in a strong magnetic field in the endcap region because of the solenoid which supplies the SDC tracking field. This flux must be returned through the endcap region of the electromagnetic (EM) and hadronic (HAD) calorimeters. Since magnetic fields are known to induce changes in the light output of plastic scintillator, the endcap will need to be recalibrated once the solenoid is energized. In addition, gradients in the field will create nonuniformities in the calorimetric response. The "induced constant term" in the EM and HAD calorimeters due to the inhomogeneities induced by the magnetic field are here evaluated in order to see that SDC performance specifications are not compromised.

2. The Magnetic Field and Plastic Scintillator

The magnetic field in SDC has been modeled [1] in 3 dimensions. The mesh for this model is sufficiently fine that gradients are easily evaluated. In Fig. 1 is shown the total field strength, in Tesla, as a function of longitudinal location (z) in the EM endcap. The radii shown span the region occupied by the endcap calorimeters. Typical longitudinal gradients are 2 kG across the EM endcap at both small and large radii. In comparison, the field strength in the HAD1 compartment ($5.0 < z < 5.9$ m) is shown in Fig. 2 as a function of z for several radii. The gradient is ~ 14 kG across the longitudinal extent of HAD1 in this case. Although the gradients are large, the SDC specification is sufficiently loose that no new provisions in hadronic calorimetry are required [2].

The fractional shift [3] in light output for SCSN38 scintillator is shown in Fig. 3. A linear representation leads to a slope of 0.6%/kG. Looking at Fig. 2, that slope means that the HAD1 compartment in the endcap sees a light output gradient of $< 6.4\%$ with a shift (to be recalibrated) of a few % [2]. For the EM region, a roughly uniform field of 12 kG leads to a shift of $\sim 8\%$, while a gradient of 2 kG causes a longitudinal inhomogeneity of 1.2%.

The transverse segmentation of the SDC EM endcap is shown in Fig. 4. The basic tile size in (η, ϕ) space is (0.05, 0.05). At large η , the EM shower would become a large fraction of this tile size, and consequently, the segmentation is increased [4]. In the region of η from 2 to 3, the field has a transverse gradient of ~ 5 kG/m. Converting from radius to η

using Fig. 4, one finds that over an EM tower, the field gradient is ~ 0.6 kG. At lower η , from 1.4 to 2 in η , the field gradient is also roughly 5 kG/m. A unit of η is physically larger, but the segmentation of a tower is finer. Consequently, the gradient over an EM tower is again ~ 0.4 to 0.6 kG/tower. Therefore, the transverse gradients in the SDC EM endcap calorimeter lead to $\sim 0.36\%$ nonuniformities in light output. Since the allowed transverse nonuniformity in a tile has already been evaluated to be $\leq 2\%$ [5], one can conclude that transverse gradients are under control. It remains to evaluate the effect of the longitudinal gradients on the energy response of the EM calorimeter.

3. EM Showers and Longitudinal Gradients

Rather than using the EGS Monte Carlo, test beam data from the "Hanging File" (HF) test calorimeter was used. The EM compartment in one incarnation of the HF stack consisted of 40 layers of $1/8$ " Pb, which is quite similar to the SDC configuration [5]. That configuration data set [6] was used for this analysis. Typical event profiles of 170 GeV electron showers are shown in Fig. 5. Layer numbers 41 to 95 correspond to 45 plates of 1" Fe. Note the quite uniform shower development of the electron showers. Fluctuations are fairly small, and are expected to consist largely of a shift of the shower shape by plus and minus about 1 radiation length. This shift consists in the smearing of a fixed shower shape by the fluctuations in the conversion point.

A model calorimeter with a 40% longitudinal gradient was "constructed" and HF electron data were thrown upon it. The resulting fractional energy shift had a mean of 24%. The distribution of this shift is shown in Fig. 6a. The rms of that distribution is 1.62%. Note that the homogeneous input HF data had a 1.5% rms at this energy. Therefore, the 40% gradient induced an uncorrected 0.61% constant term into the energy resolution. Note that the expected gradient throughout the EM compartment is expected to be $< 6.4\%$. Therefore, since the error is roughly linear in the magnitude of the gradient, the real effect is expected to be well within the SDC specifications [5].

Furthermore, the gradient causes a longitudinal inhomogeneity. Longitudinal segmentation can be used to reduce the effect of this type of inhomogeneity, as is the case for radiation induced nonuniformity. In Fig. 6b is shown the correlation between the energy at plate 9, the "shower maximum" (SM) detector and the fractional energy shift for a 40% gradient. Clearly, the effect of that gradient can be reduced by using SM information, just as it is used in reducing the induced constant term due to radiation damage.

References

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4. SDC Parameters Book, SDT-000010, September 14, 1992.
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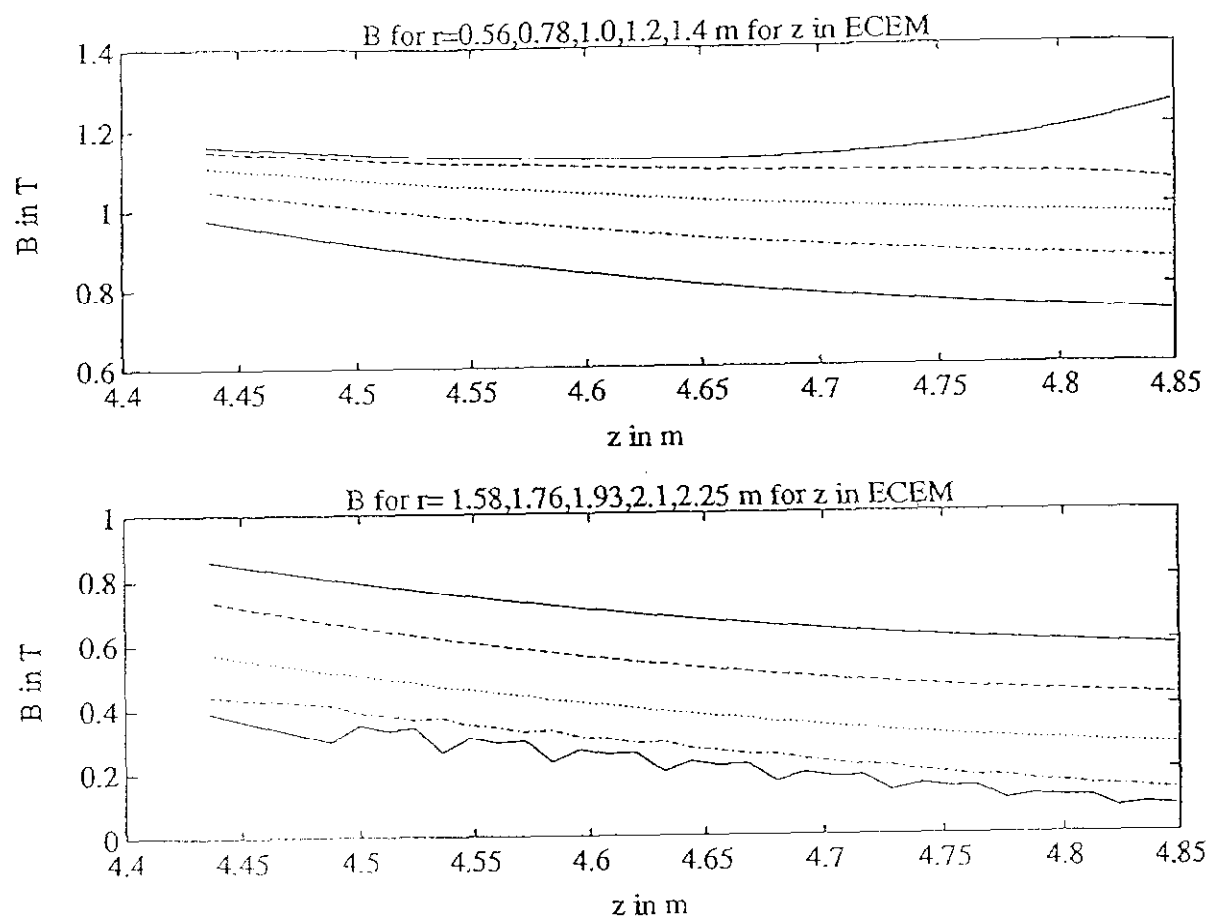


Fig. 1 Magnetic field strength as a function of depth in the SDC EM calorimeter in the endcap region.
a. radius = 0.56, 0.78, 1.0, 1.2 and 1.4 m
b. radius = 1.58, 1.76, 1.93, 2.1 and 2.25 m

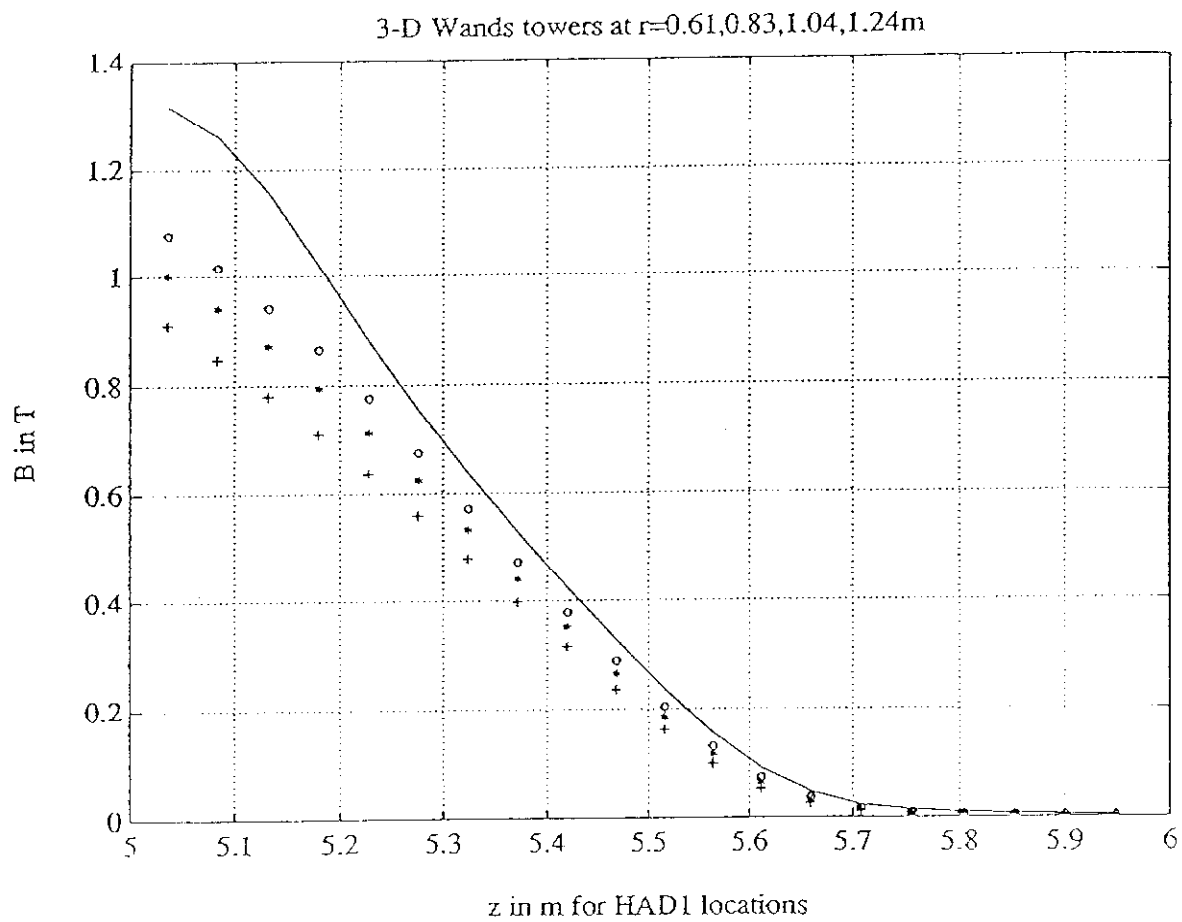


Fig. 2 Magnetic field strength as a function of depth in the SDC HAD1 calorimeter in the endcap region. The radii are $= 0.61, 0.83, 1.04$ and 1.24 m.

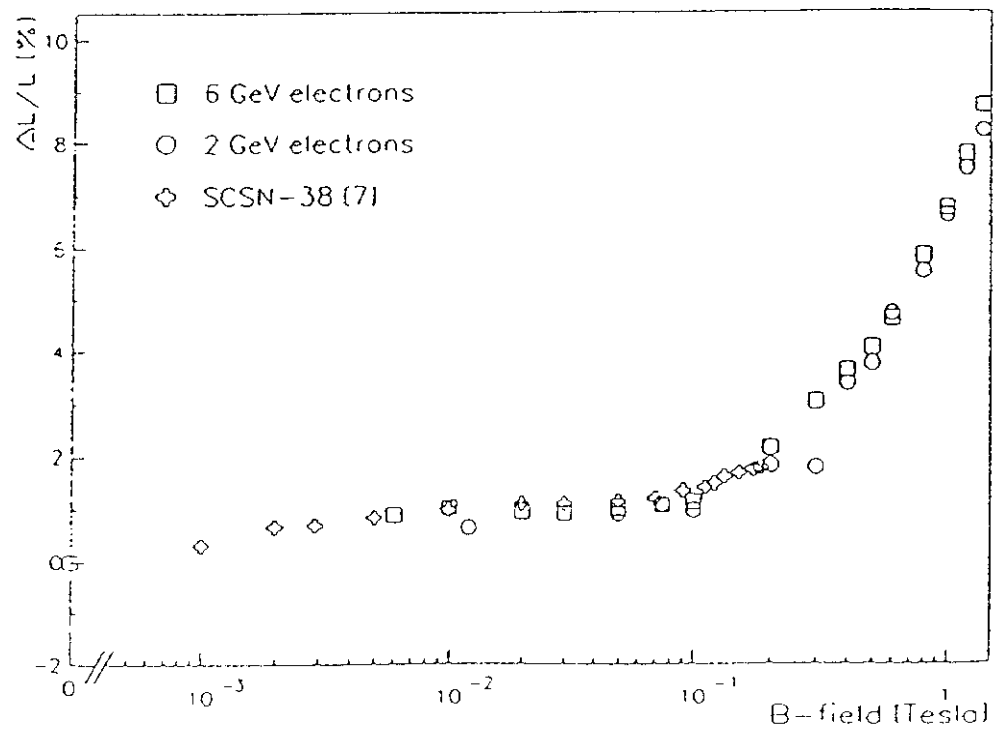


Fig. 3 Fractional light output change in SCSC38 plastic scintillator as a function of applied magnetic field strength. A linear relationship has a slope of roughly 0.6%/kG.

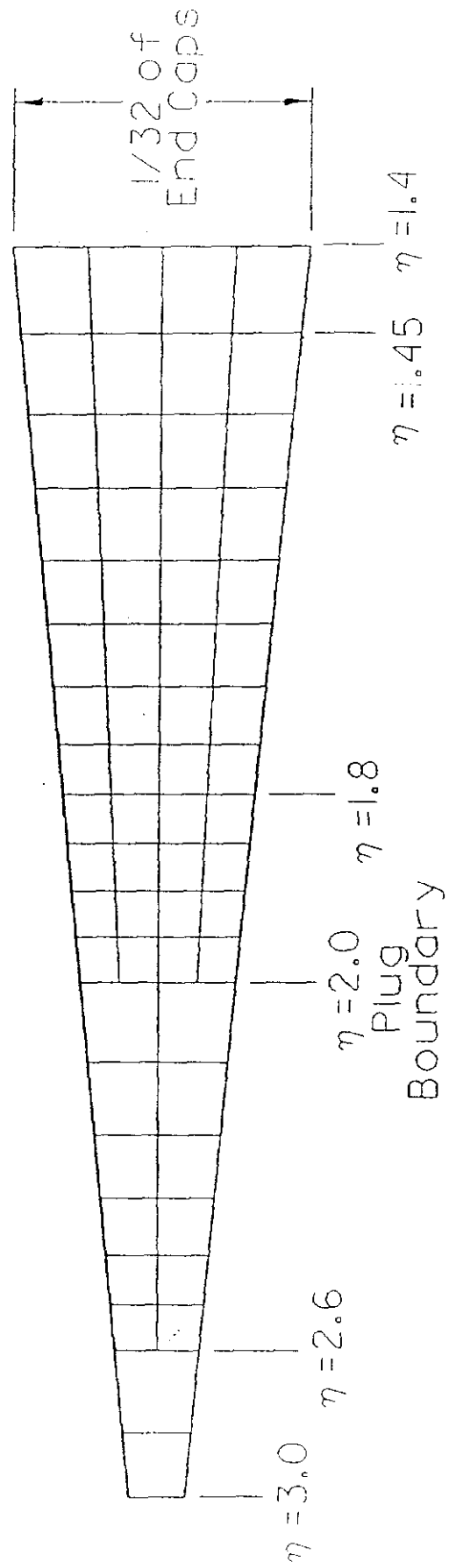


Fig. 4 Geometry of the EM segmentation in η and ϕ in the endcap. The endcap region spans the interval (1.4, 3.0) in η . The basic segmentation is 0.05×0.05 until the tile size would become comparable to the shower size, whereupon the segmentation increases to 0.1×0.1 and thence to 0.2×0.2 .

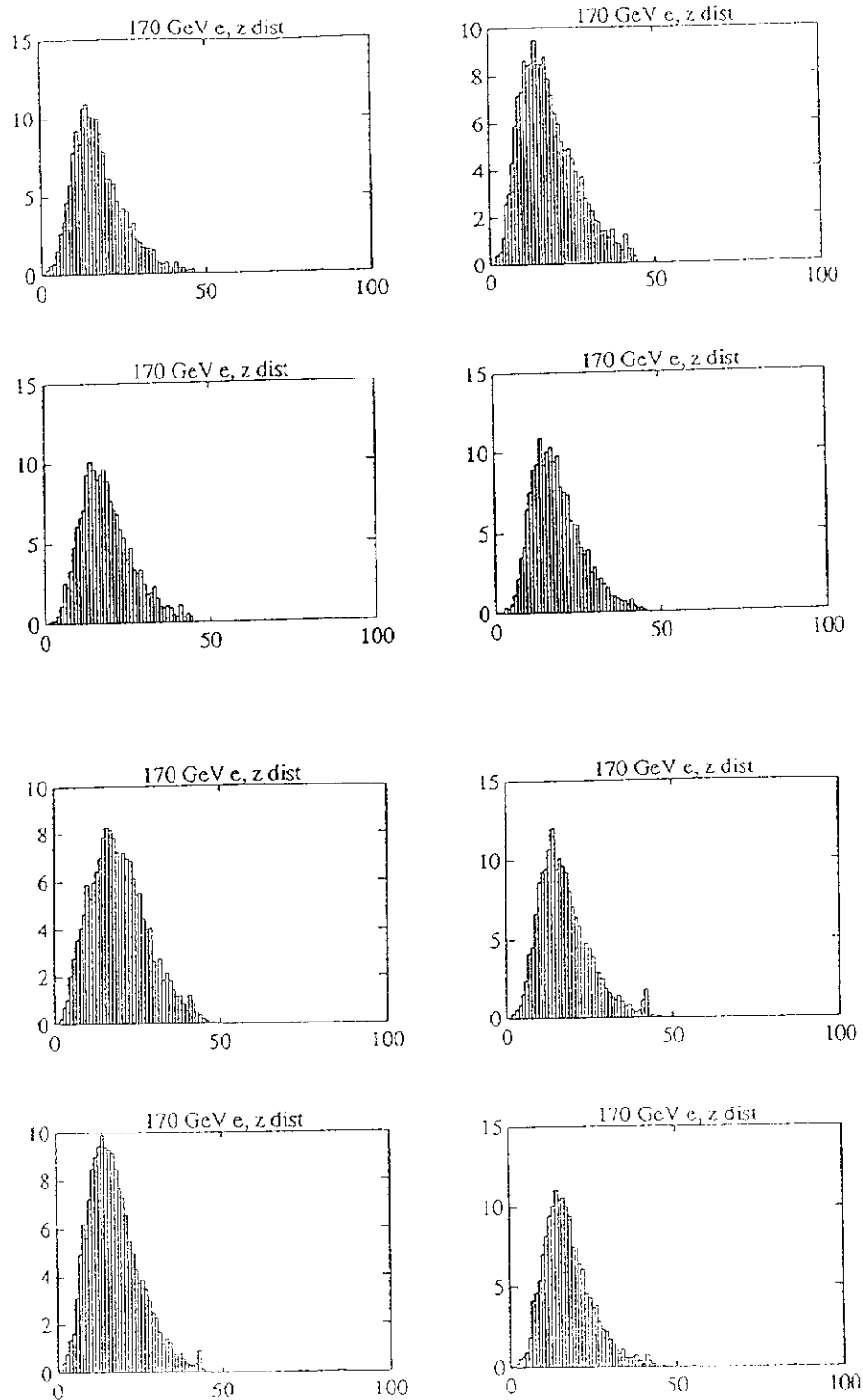


Fig. 5 Longitudinal shower distributions of 8 typical electron showers. The data is from the HF test beam set, and the stack consists of 40 plates of 1/8" Pb initially, followed by 55 plates of 1" Fe. Note the uniformity in shape of the electron showers. The mean energy for many showers is 171 GeV with a rms/mean of 1.5%.

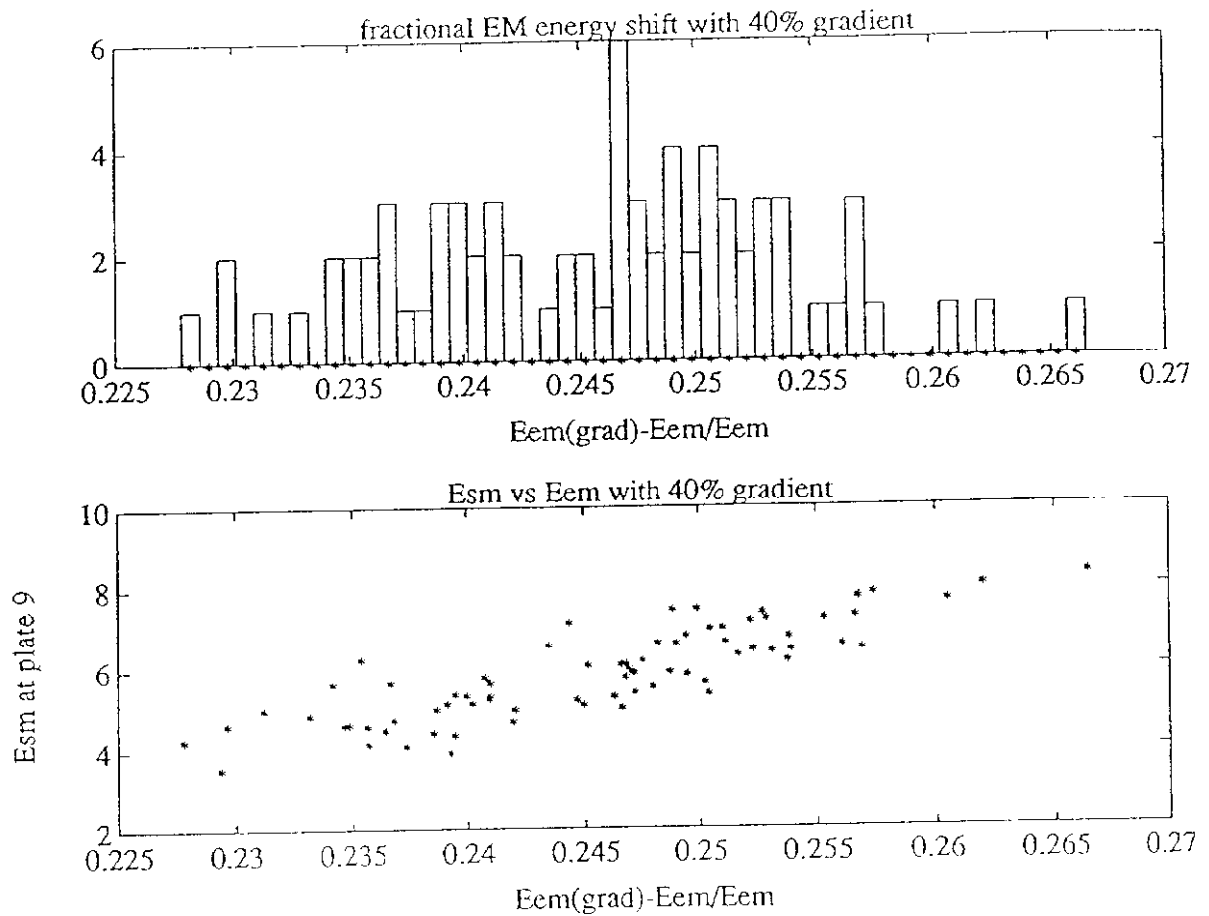


Fig. 6 a. Fractional energy shift for a 40% longitudinal gradient in the EM light output response. The mean shift is 24.5% with a rms = 1.62%. This implies an induced constant term of 0.61%.

b. Energy deposited at tile # 9 (SM) plotted against the fractional energy shift for a 40% longitudinal EM gradient. The observed correlation means that Esm information may be used to reduce the induced constant term.